

FINITE DIFFERENCE METHODS

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THE CFD APPROACH

- Assembling the governing equations
- Identifying flow domain and boundary conditions
- Geometrical discretization of flow domain
- *Discretization of the governing equations*
- *Incorporation of boundary conditions*
- Solution of resulting algebraic equations
- Post-solution analysis and reformulation, if needed

OUTLINE

- Basics of finite difference (FD) methods
- FD approximation of arbitrary accuracy
- FD formulas for higher derivatives
- Application to an elliptic problem
- FD for time-dependent problems
- FD on non-uniform meshes
- Closure

Basics of Finite Difference Methods

- FD methods are quite old and somewhat dated for CFD problems but serve as a point of departure for CFD studies
- The principal idea of CFD methods is to replace a governing partial differential equation by an *equivalent and approximate* set of algebraic equations
- Finite difference techniques are one of several options for this *discretization* of the governing equations
- In finite difference methods, each derivative of the pde is replaced by an equivalent finite difference approximation

Basics of Finite Difference Methods

- The basis of a finite difference method is the Taylor series expansion of a function.
- Consider a continuous function $f(x)$. Its value at neighbouring points can be expressed in terms of a Taylor series as

$$(1) f(x + \Delta x) = f(x) + \frac{df}{dx} (\Delta x) + \frac{d^2f}{dx^2} (\Delta x^2) / 2! + \dots + \frac{d^n f}{dx^n} (\Delta x^n) / n! + \dots$$

- The above series converges if Δx is small and $f(x)$ is differentiable
- For a converging series, successive terms are progressively smaller

FD Approximation for a First Derivative

- The terms in the Taylor series expansion can be rearranged to give

$$df/dx = [f(x+ \Delta x) - f(x)] / \Delta x - d^2f/dx^2 (\Delta x)/2! - \dots - d^n f/dx^n (\Delta x^{n-1})/n! - \dots$$

Or

$$(2) \quad df/dx \approx [f(x+ \Delta x) - f(x)] / \Delta x + O(\Delta x)$$

- Here $O(\Delta x)$ implies that the leading term in the neglected terms is of the order of Δx , i.e., the error in the approximation reduces by a factor of 2 if Δx is halved.
- Equation (2) is therefore a first order-accurate approximation for the first derivative.

Other Approximations for a First Derivative

- Other approximations are also possible. Writing the Taylor series expansion for $f(x - \Delta x)$, we have

$$(3) \quad f(x - \Delta x) = f(x) - \frac{df}{dx}(\Delta x) + \frac{d^2f}{dx^2}(\Delta x^2)/2! - \dots + \frac{d^nf}{dx^n}(\Delta x^n)/n! +$$

- Equation (3) can be rearranged to give another first order approximation :

$$(4) \quad \frac{df}{dx} \approx [f(x) - f(x - \Delta x)] / \Delta x + O(\Delta x)$$

- Subtracting (3) from (1) gives a second order approximation for df/dx :

$$(5) \quad \frac{df}{dx} \approx [f(x + \Delta x) - f(x - \Delta x)] / (2\Delta x) + O(\Delta x^2)$$

FD Approximations on a Uniform Mesh

- Consider a uniform mesh with a spacing of Δx over an interval $[0, L]$
- Denoting the mesh index by i , we can write

$$f_i = f(x_i) = f(i \Delta x) \quad \text{and} \quad f_{i+1} = f[(i+1) \Delta x] \quad \text{and so on}$$

- Then

$$(2) \Rightarrow \quad df/dx \approx [f(x+ \Delta x) - f(x)] / \Delta x \quad = (f_{i+1} - f_i)/\Delta x + O(\Delta x)$$

$$(3) \Rightarrow \quad df/dx \approx [f(x) - f(x- \Delta x)] / \Delta x \quad = (f_i - f_{i-1})/\Delta x + O(\Delta x)$$

$$(5) \Rightarrow \quad df/dx \approx [f(x+ \Delta x) - f(x-\Delta x)] / (2\Delta x) = (f_{i+1} - f_{i-1})/(2\Delta x) + O(\Delta x^2)$$

are the	forward	“one-sided”
	backward	“one-sided”
	central	“symmetric”

differencing formulas, respectively, for df/dx at x or node i

- One-sided formulas are necessary at ends of domains

Higher Order Accuracy

- Higher order of accuracy of approximation can be obtained by including more number of adjacent points
- Let us seek a third-order, one-sided approximation for $u(x)$. This requires four points and will be of the form

$$(6) \quad (du/dx)_i = [a u_i + b u_{i+1} + c u_{i+2} + d u_{i+3}] / \Delta x + O(\Delta x^3)$$

- This is equivalent to writing (6) as

$$(7) \quad (du/dx)_i = [a u_i + b u_{i+1} + c u_{i+2} + d u_{i+3}] / \Delta x + (f) d^2u/dx^2(\Delta x) \\ + (g) d^3u/dx^3(\Delta x^2) + (e) d^4u/dx^4(\Delta x^3)$$

or

$$(8) \quad a u_i + b u_{i+1} + c u_{i+2} + d u_{i+3} = + du/dx (\Delta x) + (f) d^2u/dx^2(\Delta x^2) \\ + (g) d^3u/dx^3(\Delta x^3) - (e) d^4u/dx^4(\Delta x^4)$$

- How to find a, b, c and d?

Third-Order, One-sided Formula

- Expand u_{i+1} , u_{i+2} and u_{i+3} in Taylor series about u_i :

$$(9a) \quad u_{i+1} = u(x + 1\Delta x) = u(x) + \frac{du}{dx} (\Delta x) + \frac{d^2u}{dx^2} (\Delta x)^2 / 2! + \dots$$

$$(9b) \quad u_{i+2} = u(x + 2\Delta x) = u(x) + \frac{du}{dx} (2\Delta x) + \frac{d^2u}{dx^2} (2\Delta x)^2 / 2! + \dots$$

$$(9c) \quad u_{i+3} = u(x + 3\Delta x) = u(x) + \frac{du}{dx} (3\Delta x) + \frac{d^2u}{dx^2} (3\Delta x)^2 / 2! + \dots$$

- Find $\{ a u_i + b (9a) + c(9b) + d (9c) \}$ and rearrange to get

$$(10) \quad a u_i + b u_{i+1} + c u_{i+2} + d u_{i+3} = p u + q \frac{du}{dx} (\Delta x) + r \frac{d^2u}{dx^2} (\Delta x)^2 + s \frac{d^3u}{dx^3} (\Delta x)^3 + t \frac{d^4u}{dx^4} (\Delta x)^4$$

- Compare the coefficients of (8) and (10) to get

$$a = -11/6 \quad b = 3 \quad c = -3/2 \quad d = 1/3$$

or

$$(11) \quad \left(\frac{du}{dx}\right)_i = [-11 u_i + 18 u_{i+1} - 9 u_{i+2} + 2 u_{i+3}] / (6\Delta x) + O(\Delta x^3)$$

Higher Derivatives

- Finite difference approximation for second derivative:

$$\begin{aligned}d^2u/dx^2)_i &= [d/dx(du/dx)]_i \\ &\approx [(du/dx)_{i+1/2} - (du/dx)_{i-1/2}] / \Delta x \\ &\approx [(u_{i+1} - u_i) / \Delta x - (u_i - u_{i-1}) / \Delta x] / \Delta x\end{aligned}$$

or

$$(12) \quad d^2u/dx^2)_i \approx [(u_{i+1} - 2u_i + u_{i-1})] / \Delta x^2$$

- Taylor series evaluation of equation (11) shows that the approximation is second order accurate; thus,

$$(12a) \quad d^2u/dx^2)_i = [(u_{i+1} - 2u_i + u_{i-1})] / \Delta x^2 + O(\Delta x^2)$$

- Note that use of central differences for the second derivative requires *three* points, viz., (i-1), i, (i+1), for a second order accurate formula

Other Formulas for Higher Derivatives

- Using forward differencing throughout, one can get the following first order accurate formula involving three points for the second derivative:

$$\begin{aligned} d^2u/dx^2)_i &= [d/dx(du/dx)]_i \approx [(du/dx)_{i+1} - (du/dx)_i] / \Delta x \\ &\approx [(u_{i+2} - u_{i+1}) / \Delta x - (u_{i+1} - u_i) / \Delta x] / \Delta x \end{aligned}$$

or

$$(13) \quad d^2u/dx^2)_i \approx [(u_{i+2} - 2u_{i+1} + u_i)] / \Delta x^2 + O(\Delta x)$$

- A central, second order scheme for the third derivative needs four points:

$$(14) \quad d^3u/dx^3)_i = [(u_{i+2} - 2u_{i+1} + 2u_{i-1} - u_{i-2})] / (2\Delta x^3) + O(\Delta x^2)$$

- If p = order of derivative, q = order of accuracy and n = no of points, then

$$n = p + q - 1 \quad \text{for central schemes}$$

$$n = p + q \quad \text{for one-sided schemes}$$

Mixed Derivatives

- Mixed derivatives can occur as a result of coordinate transformation to a non-orthogonal system (for example, to take account of non-regular shape of the flow domain).
- Straightforward application of the method for higher derivatives:

$$\begin{aligned}\partial^2 u / \partial x \partial y)_{i,j} &= [\partial / \partial x (\partial u / \partial y)]_{i,j} \\ &\approx [(\partial u / \partial y)_{i+1,j} - (\partial u / \partial y)_{i-1,j}] / (2\Delta x) \\ &\approx [(u_{i+1,j+1} - u_{i+1,j-1}) / 2\Delta y - (u_{i-1,j+1} - u_{i-1,j-1}) / 2\Delta y] / (2\Delta x)\end{aligned}$$

or

$$(15) \quad \partial^2 u / \partial x \partial y)_{i,j} \approx [(u_{i+1,j+1} - u_{i+1,j-1} - u_{i-1,j+1} + u_{i-1,j-1})] / (4 \Delta x \Delta y) + O(\Delta x^2, \Delta y^2)$$

- A large variety of schemes possible

Example: 2-D Poisson Equation

$$(16) \quad \partial^2 u / \partial x^2 + \partial^2 u / \partial y^2 = f \quad 0 \leq x \leq L \text{ and } 0 \leq y \leq W$$

with Dirichlet boundary conditions: $u(x,y) = g(x,y)$ on boundary

- Write $\partial^2 u / \partial x^2)_{i,j} \approx [(u_{i+1,j} - 2u_{i,j} + u_{i-1,j})] / (\Delta x^2) + O(\Delta x^2)$
and $\partial^2 u / \partial y^2)_{i,j} \approx [(u_{i,j+1} - 2u_{i,j} + u_{i,j-1})] / (\Delta y^2) + O(\Delta y^2)$
and substitute in (16) to get

$$(17) \quad [(u_{i+1,j} - 2u_{i,j} + u_{i-1,j})] / (\Delta x^2) + [(u_{i,j+1} - 2u_{i,j} + u_{i,j-1})] / (\Delta y^2) \\ = f_{ij} + O(\Delta x^2, \Delta y^2)$$

- With Dirichlet boundary conditions, equation (17) would be valid for
 $2 \leq i \leq N_i$ $2 \leq j \leq N_j$
- Results in $(N_i - 1) \times (N_j - 1)$ algebraic equations to be solved for $u(i,j)$

Poisson Equation: Other Boundary Conditions

- Normal gradient specified, e.g, $du/dy = c_1$ for all i at $j = N_j$
- Values of $u(i, N_j)$ not known and have to be determined
- For these boundary points,

$$du/dy \approx (u_{i,N_j-1} - u_{i,N_j}) / \Delta y = c_1 \quad \text{“first order accurate”}$$

or $u_{i,N_j} - u_{i,N_j-1} = c_1 * \Delta y$

- Equations for the interior points remain the same
- For second order accuracy, one can write

$$du/dy \approx (a_{i,N_j-2} + b_{i,N_j-1} + c_{i,N_j}) / \Delta y = c_1$$

- Convective boundary condition: $q''w = h*(u_{inf} - u_w)$; h and u_{inf} given
- This can be implemented by noting that $q''w = -kdu/dy$
- Thus, $h*(u_{inf} - u_{i,N_j}) = -k*(a_{i,N_j-2} + b_{i,N_j-1} + c_{i,N_j}) / \Delta y$
which gives the necessary algebraic equation for the boundary point.
- $(N_i-1) \times (N_j)$ algebraic equations to be solved for $u(i,j)$

Discretization of Time Domain

- Consider the unsteady heat conduction problem: $\partial T / \partial t = \partial^2 T / \partial x^2$ (18)

- Denote $T(x, t) = T(i \Delta x, n \Delta t) = T_i^n = T_{i,n}$

- We seek discretization of eqn. (18) of the form

$$(19) \quad \partial T / \partial t \Big|_{i,n} = \partial^2 T / \partial x^2 \Big|_{i,n}$$

- Evaluate LHS of (19) using forward differencing as

$$(20) \quad \partial T / \partial t \Big|_{i,n} = (T_{i,n+1} - T_{i,n}) / \Delta t + O(\Delta t)$$

- But several options for RHS even if we choose, say, central differencing for $\partial^2 T / \partial x^2$

Explicit and Implicit Schemes

- Put $\partial^2 T / \partial x^2 \Big|_{i,n} = (T_{i+1,n} - 2 T_{i,n} + T_{i-1,n}) / \Delta x^2 + O(\Delta x^2)$ (21)

and substitute (20) and (21) in (19) to get

- Explicit equation for $T_{i,n+1}$:

$$(T_{i,n+1} - T_{i,n}) / \Delta t = (T_{i+1,n} - 2 T_{i,n} + T_{i-1,n}) / \Delta x^2$$

or $T_{i,n+1} = T_{i,n} + \Delta t / \Delta x^2 (T_{i+1,n} - 2 T_{i,n} + T_{i-1,n}) + O(\Delta t, \Delta x^2)$ (22)

- Put $\partial^2 T / \partial x^2 \Big|_{i,n} = (T_{i+1,n+1} - 2 T_{i,n+1} + T_{i-1,n+1}) / \Delta x^2 + O(\Delta x^2)$ (23)

and substitute (20) and (21) in (19) to get

- Implicit equation for $T_{i,n+1}$:

$$(T_{i,n+1} - T_{i,n}) / \Delta t = (T_{i+1,n+1} - 2 T_{i,n+1} + T_{i-1,n+1}) / \Delta x^2 \quad \text{or}$$

$$(1 + 2 \Delta t / \Delta x^2) T_{i,n+1} = T_{i,n} + \Delta t / \Delta x^2 (T_{i+1,n+1} + T_{i-1,n+1}) + O(\Delta t, \Delta x^2) \quad (24)$$

Other Schemes

- Put $T_{i,n} = (T_{i,n-1} + T_{i,n+1})/2$ to get

- the DuFort-Frankel scheme:

$$(25) \quad (T_{i,n+1} - T_{i,n-1}) / (2\Delta t) = (T_{i+1,n} - T_{i,n+1} + T_{i,n-1} + T_{i-1,n}) / \Delta x^2 + O(\Delta t^2, \Delta x^2)$$

- Explicit, second order accurate and unconditionally stable

- Evaluate RHS at $(n+1/2)$ as $(RHS_{i,n} + RHS_{i,n+1})/2$ and put in (19) to get

- the Crank-Nicolson scheme:

$$(26) \quad -r T_{i-1,n+1} + (2+2r) T_{i,n+1} - r T_{i+1,n+1} = r T_{i-1,n} + (2-2r) T_{i,n} + r T_{i+1,n}$$

where $r = \Delta t / \Delta x^2$

- Implicit, second order accurate and unconditionally stable

Explicit vs Implicit Methods

- Explicit methods are simple to program and allow marching forward in time point by point from given initial condition
- Implicit methods are more difficult to program and require simultaneous solution of algebraic equations at each time step to get the solution
- Explicit methods are generally less stable than implicit methods and may give unphysical solutions if the marching time step is too large
- Too-large a time step, which is possible with implicit methods, may result in less accuracy and instability in non-linear problems

Non-uniform Meshes

- The methods discussed above can be extended to non-uniform meshes but may have to be done with care so as not to lose an order of accuracy :

- $$\frac{du}{dx}I = \frac{\{u_{i+1} (\Delta x_{i-1})^2 - u_{i-1} (\Delta x_i)^2 + u_i [(\Delta x_{i-1})^2 - (\Delta x_i)^2]\}}{\{(\Delta x_{i-1}) (\Delta x_i) (\Delta x_{i-1} + \Delta x_i)\}} + O\{(\Delta x_i)^2\}$$

where $\Delta x_i = x_{i+1} - x_i$ etc

- Highly non-uniform and distorted meshes should be avoided where possible

Closure

- Finite difference techniques can be used to systematically develop approximate formula for derivatives of a function on a structured mesh
- Formulas of desired accuracy level can be obtained for a derivative provided sufficient number of points are included in the formula
- Discretized equations of partial differential equations can be obtained by replacing the derivatives with finite difference formulas
- The resulting equations are coupled of algebraic equations which may be non-linear for non-linear equations and require simultaneous solution
- The issue of stability is brought into play for time-dependent problems

Thank you for your attention.

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in case of queries